LABORATORY STUDIES ON THE INFRARED ABSORPTIONS OF HYDROGENATED CARBON-SILICON CLUSTERS: DIRECTING THE IDENTIFICATION OF ORGANOMETALLIC SICH_x SPECIES TOWARD IRC +10216

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ABSTRACT

Silicon-bearing species of the molecular formula SiCH_x (x = 1-6) are plausible candidates for organometallic molecules that might be present in the circumstellar envelope of IRC +10216. Combining matrix-isolation laboratory experiments with computed vibrational spectra of these molecules, we have observed several fundamentals of the new silicon-containing radicals SiCH₃ and SiCH₅, in addition to the well-known SiCH₄ and SiCH₆ species, via infrared spectroscopy. These bands can now be used to search for hydrogen-deficient silicon-carbon–carrying molecules in the circumstellar envelope of the carbon star IRC +10216; the column densities of these species will yield valuable clues to the formation of silicon carbide molecules (SiC), together with their aggregates and partially hydrogenated clusters in the interstellar medium.

Subject headings: astrochemistry — ISM: molecules — methods: laboratory — molecular processes

1. INTRODUCTION

During the last decade, the search for interstellar carbon- and silicon-bearing molecules has received special attention in connection with the study of the atmospheres of dying carbons stars (Rittby 1992; Van Orden et al. 1994; Nielsen et al. 1997; Gordon et al. 2000; Botschwina et al. 2002). Since the identification of silane (SiH_4) in the circumstellar envelope of the carbon star IRC+10216 (CW Leo) (Barrett 1978; Goldhaber & Betz 1984), astronomers have identified 10 silicon-bearing molecules in the gas phase of cold molecular clouds, star-forming regions, and circumstellar envelopes of evolved carbons stars. These species of molecules can be classified into simple diatomic molecules of the general chemical formula SiX (group 1), silicon carbides holding the generic composition SiC_n (group 2), and silicon cyanides (group 3). The diatomic species silicon carbide (SiC), silicon sulfide (SiS), and silicon nitride (SiN) are important constituents of the circumstellar envelope of CW Leo (Cernicharo et al. 1989; Speck et al. 1997; Bieging & Nguyen 1989; Turner 1992; Wienkoop et al. 1997; Grutter et al. 1997). The second group, silicon carbides, contains three members: the cyclic $C_{2\nu}$ symmetric silicon dicarbide molecule, SiC₂ (Thaddeus et al. 1984; Sarre et al. 1996; Chandra & Sahu 1993; Gensheimer & Snyder 1997; Lloyd Evans et al. 2000); a polar, rhomboidal bicyclic SiC₃ species (McCarthy et al. 1999; Apponi et al. 1999); and a closed-shell linear chain molecule of SiC₄ (Ohishi et al. 1989). Silicon carbide and silicon dicarbide have similar column densities, whereas the higher members silicon tricarbide (SiC_3) and silicon tetracarbide (SiC4) are 200 times less abundant, with silicon tricarbide 50% more copious than silicon tetracarbide. The last group contains the SiCN and SiNC isomer pair (silicon monocyanide and silicon monoisocyanide), which have also

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been detected in the envelope of this carbon star (Guelin et al. 2000, 2004). Diatomics such as SiS were also detected in vibrationally excited ($\nu = 1$) states (Turner 1987); silicon monoxide (SiO) has only been observed in star-forming regions such as Sagittarius B2 (Peng et al. 1995; Schilke et al. 1997) and to a lesser extend in cold molecular clouds (Ziurys et al. 1989), but not in circumstellar envelopes of carbon-rich stars.

In addition to molecular lines, broad infrared emission features peaking between 11 and 11.5 μ m have been observed (Goebel et al. 1995; Mutschke et al. 1999; Papoular et al. 1998; Blanco et al. 1998; Speck et al. 1997, 1999). They have been attributed to silicon carbide grains, which are believed to be significant constituents of circumstellar envelopes of dying carbons stars; alternatively, Blanco et al. (1994) suggested amorphous carbon and silicon carbide mixtures as potential carriers. Pillinger & Russel (1993) also identified silicon carbide particles in meteorites, and their role as cometary dust particles should be pointed out as well (Orofino et al. 1994). This is remarkably distinct from oxygen-rich evolved stars, which show broad 18 μ m peaks from silicate-metal oxide dust grains (Wada et al. 2003; Greenberg 2002; Draine 2003).

Although the presence of silicon carbide in the circumstellar outflow of carbon-rich stars is well established (Clement et al. 2003), surprising little is known about the formation mechanism and chemistry of these stars. Walmsley et al. (1999) investigated the silicon chemistry in photodissociation regions (PDRs), but included predominantly silicon-oxygen chemistry. Similarly, chemical models of hot molecular cores have not coupled the silicon and carbon chemistry, but focus predominantly on silicon-oxygen and silicon-nitrogen chemical processing (MacKay 1995). The sole exceptions have been the postulated formation of silicon carbide via the reaction of silicon atoms with methyl (CH₃) and/or carbene (CH₂) followed by a stepwise photodissociation of the initial products to yield silicon carbide. Unfortunately, despite the incorporation of small hydrogenated silicon carbon clusters such as HCSi and SiCH₂, the models of Willacy & Cherchneff (1998) and Walmsley et al. (1999) could not yield conclusive evidence as to how silicon carbide grains are actually formed in circumstellar envelopes. This suggests that crucial formation routes are lacking. From the chemical viewpoint, it is quite surprising that with the exception of HCSi and SiCH₂, hydrogenated binuclear carbon-silicon clusters of the generic formula $SiCH_x$ (x = 6-3) have neither been incorporated into chemical models of, for instance, the carbon star IRC +10216, nor searched for in extraterrestrial environments. In the circumstellar envelope, a rich silane and methane chemistry can be driven by a photolysis of the silane and methane precursor molecule, thus forming the methyl (CH₃) and silyl (SiH₃) radicals. The latter could recombine to form the fully hydrogenated, binuclear carbon-silicon cluster silvlmethane, SiCH₆, in a three-body reaction; if the reaction zone is too far from the central star, the SiCH₆ molecule cannot be stabilized by a threebody collision and decomposed to SiCH₅ isomers (recall that the number density drops with the inverse square of the radius from the centrals star). Successive photolysis of these species can yield silicon carbide (SiC) as a final product via stepwise atomic and/or molecular hydrogen elimination.

To test the validity of this hypothesis, it is important to search for SiCH_x (x = 6-3) species in circumstellar envelopes of carbons stars. One drawback of a potential astronomical observation is the relatively small dipole moments of these molecules, less than 1.5 debye (Bennett et al. 2005a, 2005b). Therefore, an infrared spectroscopic search of the SiCH_x species might present a viable alternative. Unfortunately, their infrared spectra—with the exception of silylmethane (CH₃SiH₃)—are elusive. In this paper, we present a matrix-isolation study of the infrared spectra of various SiCH_x isomers (x = 6-3). This will assist an astronomical search for hydrogenated carbon-silicon clusters in the circumstellar envelope of IRC +10216 and provides valuable information on potential intermediates to finally synthesize silicon carbide in the interstellar medium.

2. EXPERIMENTAL PROCEDURE

Our experiments were carried out in a contamination-free ultrahigh vacuum (UHV) matrix-isolation chamber that can be evacuated down to 5×10^{-11} torr by a magnetically suspended turbopump backed by an oil-free scroll pump. A rotatable, two-



FIG. 1.—Infrared spectrum of the silane-methane sample at 10 K.

 TABLE 1

 Infrared Absorptions of the Methane-Silane Matrix

Frequency (cm ⁻¹)	Assignment	Carrier
4523	$\nu_2 + \nu_3$	CH ₄
4364	$2 \nu_3$	SiH_4
4358	$\nu_3 + \nu_4$	CH_4
4298	$\nu_3 + \nu_4$	CH_4
4292	$\nu_1 + \nu_3$	SiH_4
4200	$\nu_1 + \nu_4$	CH_4
4112	$\nu_2 + 2\nu_4$	CH_4
3852/3895	$3 \nu_4$	CH_4
	$\nu_3 + \nu_4 + \beta$	SiH_4
3139	$\nu_2 + \nu_3$	SiH_4
	$\nu_3 + \nu_4 + \alpha$	SiH_4
3074	$\nu_3 + \nu_4$	SiH_4
3011	ν_3	CH_4
2815	$\nu_2 + \nu_4$	CH_4
2595	$2\nu_4$	CH_4
2306	$\nu_3 + \gamma$	SiH_4
	$\nu_3 + \beta$	SiH_4
2193	$\nu_3 + \alpha$	SiH_4
2164	ν_3	SiH_4
1878	$\nu_2 + \nu_4 + \alpha$	SiH_4
1300	$ u_4$	CH_4
1032	$ u_2 + eta$ / $ u_4 + \gamma$	SiH_4
1032	$ u_2 + eta$ / $ u_4 + \gamma$	SiH_4
961	ν_2	SiH_4
930	$\nu_4 + \beta$	SiH_4
914	$\nu_4 + \alpha$	SiH_4
883	$ u_4$	SiH_4

NOTE.—Here α , β , and γ denote lattice modes.

stage closed-cycle helium refrigerator is attached to the lid of the machine and holds a polished silver single crystal. The latter is cooled to 10.2 ± 0.3 K and serves as a substrate for the silane-methane ices. The ices were prepared at 10 K by depositing silane (99.99%) and methane (99.99%) at pressures of 8×10^{-8} torr for 20 minutes onto the cooled silver crystal. Figure 1 shows a typical infrared spectrum of the frost at 10 K; the absorptions are summarized in Table 1. To determine the ice thickness quantitatively, we integrated the infrared absorption features at 2164 and 883 cm^{-1} (silane) as well as 4200, 3852, and 2815 cm^{-1} (methane). The ice thicknesses are then calculated via the Lambert-Beers relationship. Considering the integrated absorption coefficients of these fundamentals, i.e., 2.5×10^{-17} and 2.0×10^{-17} cm (silane) as well as 1.6×10^{-18} , 2.0×10^{-19} , and 2.3×10^{-18} cm (methane), this gives column densities of $2.0 \pm 0.1 \times 10^{17}$ and $1.5 \pm 0.2 \times 10^{17}$ cm², respectively. Taking the densities of the silane and methane ices to be 0.77 \pm 0.03 and 0.53 \pm 0.02 g cm⁻³ (Sears & Morrison 1975), this translates into thicknesses of 140 ± 15 nm silane and 70 \pm 20 nm methane. The ices were irradiated at 10 K with 5 keV electrons generated in an electron gun at beam currents of 100 and 1000 nA by scanning the electron beam over an area of 3.0 ± 0.4 cm². Here the electron beam is used to cleave the carbon-hydrogen and silicon-hydrogen bonds of the methane and silane molecules and hence to initiate the formation of reactive radical precursors. Accounting for irradiation times of 60 minutes and the extraction efficiency of 78.8% of the electrons, this exposes the targets to up to 1.8×10^{16} electrons. To guarantee an identification of the reaction products in the ices and those subliming into the gas phase on line and in situ, a

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Fourier transform infrared spectrometer (FTIR; solid state) and a quadrupole mass spectrometer (QMS; gas phase) were used. The Nicolet 510 DX FTIR spectrometer operates in an absorption-reflection-absorption mode (reflection angle $\alpha = 75^{\circ}$; resolution $0.5-2 \text{ cm}^{-1}$; spectra were averaged for 180 s). The infrared beam is coupled via a mirror flipper outside the spectrometer, passes through a differentially pumped potassium bromide (KBr) window, is attenuated in the ice sample before and after reflection at a polished silver waver, and exits the main chamber through a second differentially pumped KBr window before being monitored via a liquid-nitrogen–cooled detector. The gas phase is sampled by a quadrupole mass spectrometer (Balzer QMG 420; 1–200 amu mass range) with electron impact ionization of the neutral molecules.

3. RESULTS

Infrared spectroscopic studies of the irradiation phase suggest that the reaction of the silane-methane matrix to electron bombardment is dictated by the initial formation of the silyl, SiH₃ (X^2A_1) , and methyl, CH₃ (X^2A_2'') , radicals plus atomic hydrogen (reactions [1] and [2])

$$CH_4(X^1A_1) \to CH_3(X^2A_2'') + H(^2S_{1/2}),$$
 (1)

$$\operatorname{SiH}_4(X^1A_1) \to \operatorname{SiH}_3(X^2A_1) + \operatorname{H}({}^2S_{1/2}).$$
 (2)

The ν_2 vibrations of both the silyl and methyl radicals are visible at 723 and 609 cm⁻¹, respectively. These data agree very well with previous matrix studies (Andrews & Wang 2002). Very strong features of well-known ethane [C₂H₆(X¹A_{1g})] and disilane [Si₂H₆(X¹A_{1g})] molecules—typical recombination products of methyl (Bennett et al. 2005a, 2005b) and silyl radicals (Sillars et al. 2004) as observed in pure methane and silane ices showed up as well (Table 2; reactions [3] and [4]):

$$CH_3(X^2A_2'') + CH_3(X^2A_2'') \to C_2H_6(X^1A_{1g}),$$
 (3)

$$SiH_3(X^2A_1) + SiH_3(X^2A_1) \to Si_2H_6(X^1A_{1g}).$$
 (4)

With increasing irradiation time, the ethane molecules are radiolyzed stepwise via the ethyl radical $[C_2H_5(X^2A')]$, ethylene $[C_2H_4(X^1A_g)]$, and the vinyl radical $[C_2H_3(X^2A')]$ to finally form acetylene $[C_2H_2(X^1\Sigma_g^+)]$; disilane was found to degrade via the disilyl radical $[Si_2H_5(X^2A')]$, two Si_2H_4 isomers, silylsilylene $[H_3SiSiH(X^1A']$, disilene $[H_2SiSiH_2(X^1A_g)]$, and disilenyl $[H_2SiSiH(X^2A)]$, to yield the bridged disilicondihydride $[SiH_2Si(X^1A_1)]$. These reaction sequences were also observed in neat methane and silane matrices at 10 K (Sillars et al. 2004, 2005; Kaiser & Osamura 2005). In addition to the absorption present in the methane and silane matrices, we also detected six prominent features of the well-known methylsilane molecule, $CH_3SiH_3(X^1A_1;$ see Table 2), which is formed via the recombination of a methyl with a silyl radical (reaction [5]):

$$SiH_3(X^2A_1) + CH_3(X^2A_2'') \to CH_3SiH_3(X^1A_1).$$
 (5)

We also identify two new SiCH₅ isomers: methylsilyl, CH₃SiH₂(X^2A'), and silylmethyl, SiH₃CH₂(X^2A'), via their absorptions at 1251 cm⁻¹ (ν_5), 1414 cm⁻¹ (ν_{12}), and 653 cm⁻¹ (ν_8), as well as 645 cm⁻¹ (ν_8) (see Figs. 2 and 3). These findings are in excellent agreement with computed frequencies of both isomers (Bennett et al. 2005a, 2005b).

TABLE 2 Compilation of Newly Observed Species in the Methane-Silane Matrix upon Electron Irradiation

Frequency				
(cm^{-1})	Fundamental	Carrier		
Hydrogenated Carbon Species				
609	ν_2	CH ₃		
826	ν_{12}	C_2H_6		
1380	ν_{16}	C_2H_6		
1462	ν_{11}	C ₂ H ₆		
2880	ν_5	C_2H_6		
2973	ν_{10}	C ₂ H ₆		
530	ν_0	C ₂ H ₅		
2848	ν_{2}	C ₂ H ₅		
2939	2 3 1/2	CaHe		
1440	V12	C ₂ H ₃		
960	V 12	C ₂ H ₄		
826	ν_{7}	$C_2\Pi_4$		
820	ν_{10}			
699	ν_8	$C_2 \Pi_3$		
2282	ν_7	$C_2 \Pi_3$		
726	ν_3	$C_2\Pi_2$		
/30	ν_5	C_2H_2		
Hydrogenated	Silicon Species			
722	ν_2	SiH ₃		
826	ν_6	Si ₂ H ₆		
934	ν_{11}	Si ₂ H ₆		
2142	$\nu_{5/4}$	Si ₂ H ₆		
852	ν_6	Si ₂ H ₅		
856	ν5	H ₃ SiSiH		
897	ν_{11}	H ₂ SiSiH ₂		
933	ν_{Λ}	H ₂ SiSiH		
Hudrogonated Carl	on Silicon Spacing	112515111		
	Join-Silicoli Species			
691	ν_5	CH ₃ SiH ₃		
863	ν_{11}	CH ₃ SiH ₃		
946	ν_{10}	CH ₃ SiH ₃		
976	$ u_4 $	CH ₃ SiH ₃		
1254	ν_3	CH ₃ SiH ₃		
2142	ν_2	CH ₃ SiH ₃		
654	ν_8	CH ₃ SiH ₂		
1414	ν_{12}	CH ₃ SiH ₂		
1251	ν_5	CH ₃ SiH ₂		
645	ν_8	CH ₂ SiH ₃		
740	ν_7	H ₂ CSiH ₂		
1947	ν_3	CH ₃ SiH		
1211	ν,	H ₃ CSi		
1387	ν ₂	H ₂ CSi		
717	- 3	H		
688	V 8	H ₂ CSi		
~ ~ ~ ~	- 4	/		



FIG. 2.—Structures of the hydrogenated silicon-carbon clusters. Bond lengths and bond angles are in Å and degrees (*dark gray*: silicon; *light gray*: carbon; *white*: hydrogen atoms). The computed frequencies and computational methods are listed in the Appendix.



Fig. 3.—Absorption features of the newly silicon-carbon bearing species $\rm SiCH_5$ and $\rm SiCH_3$ in the silane-methane matrices.

The assignment of the methylsilyl and silvlmethyl radicals raises the question whether, in a manner similar to the radiolysis of ethane and disilane, an enhanced electron exposure of the silane-methane matrices will lead to the formation of hydrogendepleted silicon-carbon clusters of the formula $CSiH_n$ (n = 4, 3, 2). Indeed, with increasing irradiation time, new absorption features appear at 742 and 1947 cm⁻¹. These absorptions could not be assigned to any of the molecules discussed above. Instead, a comparison with computed vibration frequencies of the SiCH₄ isomers (Bennett et al. 2005a, 2005b) suggest the formation of the most thermodynamically stable silene molecule $[H_2CSiH_2(X^1A_1)]$ (ν_7) and the less stable methylsilylene isomer $[CH_3SiH(X^1A')]$ (ν_3). An enhanced irradiation times lead to the synthesis of $H_3CSi(X^2A'')$ and $H_2C=SiH(X^2A')$. Here, absorptions at 1226 cm⁻¹ (ν_4) and 1371 cm⁻¹ (ν_3) were assigned to the H₃CSi(X^2A'') structure; the peak at 822 cm⁻¹ (ν_6) could be attributed to the H₂C=SiH(X^2A') isomer (Figs. 2 and 3, Table 2). These data are in excellent agreement with previous calculations (Bennett et al. 2005a, 2005b).

4. ASTROPHYSICAL IMPLICATIONS

Our studies provide reliable infrared absorptions of astrophysically important hydrogenated silicon-bearing molecules. Two classes of molecules can be identified from the present laboratory study. These are XH_3 species (X = C, Si) and XYH_r (X = C, Si; Y = C, Si; x = 1-6). The silicon-carbon-bearing dinuclear clusters SiCH_x are of great interest in the potential formation of silicon carbide molecules and their aggregates in circumstellar envelopes of dying carbon stars such as of IRC +10216. Here, the fundamentals as observed in our experiments are compiled in Table 2 and serve as valuable guidelines for prospective astronomical identifications of these molecules. We would like to stress that the actual photodissociation of silane molecules in circumstellar envelopes is expected to preferentially yield the SiH₂($X^{1}A_{1}$) and SiH($X^{2}\Pi$) species. However, these radicals can react with ubiquitous methane molecules to form methylsilane and silylmethyl/methylsilyl isomers via insertion processes followed by isomerization:

$$\operatorname{SiH}_2(X^1A_1) + \operatorname{CH}_4(X^1A_1) \to \operatorname{CH}_3\operatorname{SiH}_3(X^1A_1), \quad (6)$$

$$\operatorname{SiH}(X^{2}\Pi) + \operatorname{CH}_{4}(X^{1}A_{1}) \to \operatorname{SiH}_{2}\operatorname{CH}_{3}(X^{2}A'), \qquad (7)$$

$$SiH_3CH_2(X^2A') \to CH_3SiH_2(X^2A').$$
(8)

Considering the homologous C_2H_x (x = 1-6) series, we would like to point out that silicon-bearing species have not been

TABLE 3 Dipole Moments and Rotational Constants for SiCH₆, SiCH₅, SiCH₄, and SiCH₃ Calculated with the B3LYP/6-311G(d, p) Theory

Species	Dipole Moment (debye)	Rotational Constants (GHz)		
CH ₃ SiH ₃	0.738	56.119	10.821	10.821
CH ₃ SiH ₂	0.806	73.007	11.583	11.027
CH ₂ SiH ₃	0.650	67.012	11.952	11.475
CH ₂ SiH ₂	0.600	105.056	14.763	12.944
CH ₃ SiH	1.009	93.354	12.224	11.571
CH ₂ SiH	0.554	157.285	15.473	14.087
HCSiH ₂	0.635	161.999	15.801	14.397

detected in either the solar system or the interstellar medium so far. On the other hand, members of the homologous C_2H_x (x = 1-6) series are very abundant in our solar system and in the interstellar medium. Ethane (C_2H_6) and acetylene (C_2H_2) are actually present in the atmospheres of the giant planets Jupiter and Saturn (Fouchet et al. 1999; Courtin et al. 1984), as well as Uranus (Atreya et al. 1999; Bishop et al. 1990) and Neptune (Conrath et al. 1989; Orton et al. 1992). Ethylene (C_2H_4) was only observed in Jupiter and Neptune (Encrenaz et al. 2002; Fegley 1995). Ethane has been also identified in comets such as 153P/Ikeya-Zhang (Kawakita et al. 2003). Note that ethylene (Goldhaber et al. 1987), acetylene (Cernicharo et al. 1999), and the ethinyl radical (Fuente at al. 1998) have been detected in the circumstellar envelope of IRC +10216. We also computed the rotational constants and the dipole moments of the organosilicon species detected in our experiments (Table 3). Here the dipole moments are very small (only up to about 1 debye), so these molecules might be difficult to detect with microwave spectroscopy. Therefore, our laboratory investigations should be helpful in elucidating the hitherto poorly understood organosilicon chemistry in circumstellar envelopes of carbon-rich dying stars and also in prospective astronomical searches in the infrared regime using the NASA Infrared Telescope Facility and the Spitzer Space Telescope.

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APPENDIX

COMPUTATIONAL METHOD

The vibrational frequencies and the molecular structures of the $CSiH_x$ (x = 1-6) isomers shown in Table 4 were computed in terms of ab initio molecular orbital methods. The geometries were optimized with the hybrid density functional B3LYP method, i.e., Becke's three-parameter nonlocal exchange functional (Becke 1993) with the nonlocal correlation functional of Lee et al. (1988) and the 6-311G(d, p) basis set (Krishnan et al. 1980). To assist spectroscopic identification in circumstellar envelopes, we also obtained vibrational frequencies and infrared intensities. The coupled cluster CCSD(T) calculations (Cizek 1969; Pople et al. 1987) with the aug-cc-pVTZ basis set (Kendall et al. 1992) were also performed for the optimized structures obtained with the B3LYP method. All computations were carried out using the GAUSSIAN 98 program package (Frisch et al. 2002).

TABLE 4 UNSCALED VIBRATIONAL FREQUENCIES (cm⁻¹) OF EXPERIMENTALLY OBSERVED Hydrogenated Carbon-Silicon Clusters Calculated with B3LYP/6-311G(d, p) Theory

Mode	Frequency	Characterization		
	CH ₃ SiH ₃			
$\nu_1 a_1$	3031	CH ₃ sym. stretch		
$\nu_2 a_1$	2221	SiH ₃ sym. stretch		
$\nu_3 a_1$	1305	CH ₃ umbrella		
$\nu_4 a_1$	949	SiH ₃ umbrella		
$\nu_5 a_1$	690	C—Si stretch		
$\nu_6 a_2$	194	torsion		
ν ₇ ε	3105	CH ₃ asym. stretch		
ν ₈ ε	2221	SiH ₃ asym. stretch		
Vo 6	1469	CH ₂ deformation		
V10 C	965	SiH ₃ deformation		
V11 C	896	CH ₂ , SiH ₂ rocking		
$\nu_{12} \ e$	525	CH ₃ , SiH ₃ rocking		
	CH ₃ SiH ₂			
ν ₁ a'	3091	CH ₃ asym. stretch		
$\nu_2 a'$	3021	CH ₃ sym. stretch		
$\nu_3 a'$	2178	SiH ₂ sym. stretch		
$\nu_4 a'$	1463	CH ₃ deformation		
$\nu_5 a'$	1291	CH ₃ umbrella		
$\nu_6 a'$	938	SiH ₂ bending		
$\nu_7 a'$	851	CH ₃ , SiH ₂ rocking		
$\nu_8 a'$	681	C-Si stretch		
$\nu_9 a'$	570	CH ₃ , SiH ₂ rocking		
$\nu_{10} a''$	3112	CH_3 asym. stretch		
$\nu_{11} a''$	2203	SiH ₂ asym. stretch		
$\nu_{12} a''$	1458	CH ₃ deformation		
$\nu_{12} a''$	895	CH_2 , SiH ₂ deformation		
$\nu_{14} a''$	530	CH_2 , SiH ₂ deformation		
$\nu_{15} a''$	177	torsion		
-	CH ₂ SiH ₃			
$\nu_1 a'$	3122	CH ₂ sym. stretch		
$\nu_2 a'$	2221	SiH ₃ sym. stretch		
$\nu_{2} q'$	2185	SiH ₂ asym. stretch		
$\nu_A q'$	1405	CH ₂ bending		
$\nu_{\epsilon} a'$	951	SiH ₂ umbrella		
<i>v</i> _c <i>a</i> [']	948	SiH ₂ deformation		
$\nu_6 a'$	748	C-Si stretching		
$\nu_{0} a'$	666	SiH ₂ rocking		
$\nu_{8} a'$	525	CH, umbrella		
<i>vyu</i>	223	CH, asym strotch		
$\nu_{10} \alpha$	2214	SiH, asym. strotch		
$\nu_{11} a$	061	Sills deformation		
$\nu_{12} u$	201 9 <i>47</i>	CH. Sill, realing		
$\nu_{13} a \dots $	047 526	CH_2 , SH_3 focking		
$\nu_{14} a \dots \nu_{15} a''$	17	torsion $C\Pi_2$, SI Π_3 rocking		
	22	CH. sum attatah		
ν ₁ α ₁	5145 2245	CH ₂ sylli. stretch		
$\nu_2 a_1$	2205	SIH_2 sym. stretch		
$\nu_3 a_1$	1401	CH_2 scissor		
$\nu_4 a_1$	997	C-Si stretch		
$\nu_5 a_1$	943	SiH ₂ scissor		
$\nu_6 a_2$	730	twist		
$\nu_7 b_1$	782	CH ₂ out-of-plane		
$\nu_8 b_1$	455	SiH ₂ out-of-plane		
$\nu_9 b_2$	3231	CH ₂ asym. stretch		
$\nu_{10} \ b_2$	2288	SiH ₂ asym. stretch		
$\nu_{11} \ b_2$	841	CH ₂ rocking		
$\nu_{12} \ b_2$	485	SiH ₂ rocking		

TABLE 4—Continued

Mode	Frequency	Characterization		
H ₃ C-SiH				
<i>ν</i> ₁ <i>a</i> ′	3104	CH ₃ asym. stretch		
$\nu_2 a'$	2993	CH ₃ sym. stretch		
$\nu_3 a'$	2005	SiH stretch		
$\nu_4 a'$	1437	CH ₃ deformation		
$\nu_5 a'$	1267	CH ₃ umbrella		
$\nu_6 a'$	952	SiH bend, CH ₃ rock		
$\nu_7 a'$	662	C-Si stretch		
$\nu_8 a'$	629	SiH bend		
$\nu_9 a''$	3045	CH_3 asym. stretch		
$\nu_{10} a''$	1453	CH_3 deformation		
$\nu_{11} a^{*}$	578	CH ₃ rocking		
$\nu_{12} a$	150	torsion		
H ₃ C–Si				
<i>ν</i> ₁ <i>a</i> ′	3040	CH ₃ asym. stretch		
$\nu_2 a'$	2965	CH ₃ sym. stretch		
$\nu_3 a'$	1445	CH ₃ deformation		
$\nu_4 a'$	1259	CH ₃ umbrella		
$\nu_5 a'$	665	C-Si stretch		
$\nu_6 a'$	528	CH ₃ rocking		
$\nu_7 a''$	3072	CH ₃ asym. stretch		
$\nu_8 a''$	1359	CH ₃ deformation		
$\nu_9 a''$	598	CH ₃ deformation		
	H ₂ C=SiH			
$\nu_1 a'$	3223	CH ₂ asym. stretch		
$\nu_2 a'$	3116	CH_2 sym. stretch		
$\nu_3 a'$	2128	SiH stretch		
$\nu_4 a'$	1368	CH ₂ scissor		
$\nu_5 a'$	964	C-Si stretch		
$\nu_6 a'$	839	SiH bend, CH2 rock		
$\nu_7 a'$	521	SiH bend, CH2 rock		
$\nu_8 a''$	749	out-of-plane		
$\nu_9 a''$	256	torsion		
HC=SiH ₂				
ν ₁ a ₁	3343	CH stretch		
$\nu_2 a_1$	2254	SiH ₂ sym. stretch		
$\nu_3 a_1$	1072	C-Si stretch		
$\nu_4 a_1$	958	SiH ₂ scissor		
$\nu_5 b_1$	667	out-of-plane		
$\nu_6 b_1$	319	out-of-plane		
$\nu_7 b_2$	2265	SiH ₂ asym. stretch		
$\nu_8 b_2$	606	SiH ₂ rocking		
$\nu_9 b_2$	196	CH bend		

 $\operatorname{Note.}-Frequencies$ must to be scaled by a factor of $0.98{-}0.97$ to be compared with experimental data.

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